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Effect of shrouding CH₄ flow rate on flow field and stirring ability of coherent jet in steelmaking process

Fuhai Liu¹, Dongbai Sun¹, Rong Zhu^{2*}, Rongfang Su³ and Xueyi Wang⁴

*Correspondence:

zhurong1206@126.com

² Beijing Key Laboratory of Research Center of Special Melting and Preparation of High-end Metal Materials, University of Science and Technology Beijing, Beijing 100083, China

Full list of author information is available at the end of the article

Abstract

Characteristics of flow field and stirring ability of coherent jet with various shrouding CH₄ flow rates on the molten bath were studied by combustion experiment and numerical simulation. The axial velocity and total temperature distributions of coherent jet under hot (1700 K) and cold (298 K) ambient condition were analyzed. The Eddy Dissipation Concept model was used in simulation with detail chemical kinetic mechanisms, and the numerical simulation results were agreed well with the combustion experiment in this research. Based on the simulation and experiment results, when the CH₄ rate was 230, 207 and 184 Nm³/h, their disparity rate of average velocity and total temperature was small than 5 and 6 %, respectively, at high ambient temperature. Hence, the same stirring effect might be achieved by those three kinds of CH₄ flow rates in EAF steelmaking process. According to the industrial application research, the best CH₄ flow rate is 184 Nm³/h, which could stir molten bath well and reduce energy consumption in steelmaking process.

Keywords: Electric arc furnace, Coherent jet, Combustion experiment, Numerical simulation

Background

During the steelmaking process, the main equipment to supply oxygen into the molten bath is oxygen lance, which is used widely in the basic oxygen furnace (BOF) and the electric arc furnace (EAF). Moreover, the oxygen lance also have a key function in dephosphorization and stirring the molten pool. In order to achieve a great stirring ability and high reaction rate, the Laval nozzle is adopted to increase the velocity of oxygen jet (Deo and Boom 1993; Naito et al. 2000). In the Laval nozzle, the high pressure energy of oxygen jet is transformed into its kinetic energy, and the jet would be accelerated to 2.0 Mach number. After that, the velocity of oxygen jet begins to reduce due to entrainment of ambient gas, which forms potential core, supersonic core and subsonic zone. With increasing distance from the tip of nozzle, the impact force of the oxygen jet decreases which would reduce the mass transfer processes and the reaction rates in the furnace (Tago and Higuchi 2003; Wang et al. 2010).

To solve that problem, the coherent jet technology has been applied to electric arc furnace steelmaking process (Liu et al. 2005; Mathur 1999a), which could prolong the

length of potential core and increase the kinetic energy of the main oxygen jet. The key of this technology is suppressing the ambient gas to interact with main oxygen jet by a shrouding flame. Therefore, the oxygen jet could keep original diameter and velocity over long distance, and remain the stirring ability. Moreover, with the greater penetration capacity, the supersonic oxygen jet could deliver greater amounts of oxygen into the molten bath and decrease the splash of liquid slag, comparing with the traditional supersonic oxygen lance (Nordquist et al. (2006); Mahoney (2008)).

The coherent jet technology was proposed by Mathur and Anderson et al. (1998), and then Mathur (1999b) showed the fundamental results in industrial production. Sarma et al. (1998) analyzed the characteristics of supersonic jet and coherent jet at different ambient temperatures by combustion experiment. Alam et al. (2010) reported the behaviors of the supersonic jet with and without shrouding flame by the numerical simulation and validated their results against previously experimental data. To date, although substantial works have been performed for researching the differences between conventional oxygen jet and coherent jet (Anderson et al. 1998; Mathur 1999b; Sarma et al. 1998; Alam et al. 2010a; Mahoney 2010; Meidani et al. 2004; Liu et al. 2016; Jeong et al. 2004), little research has been paid to how the flow field is influenced by flow rate of shrouding flow.

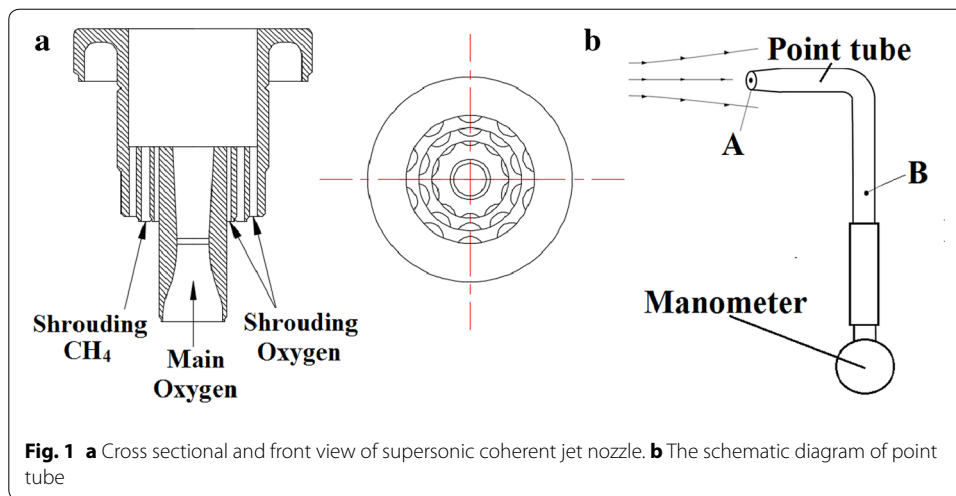
The present study is to make some contributions to address this deficit by combustion experiment and numerical simulation. During the combustion process, the process products are formed such as CH_5 , CH_2 and OH , which could not be tracked by conventional one-step complete combustion. Therefore, the numerical simulation is performed with detail chemical kinetic mechanisms to analyze the effect of process products on coherent jet flow field, and EDC model is used to analyze how the shrouding flow rate affects the potential core length and stirring ability of coherent jet under two kinds of ambient temperatures. Based on results, the metallurgical effects of various CH_4 flow rates are analyzed in a 75 t electric arc furnace.

Combustion experiment

The present research used a combustion system to produce a hot ambient temperature for the coherent jet. The details of the system were described in Ref. (Mardani et al. 2010), and only a brief description was given here. In this paper, the coherent jet lance had a water-cooled system, and the design flow rate of main oxygen was $2300 \text{ Nm}^3/\text{h}$. Moreover, the throat and exit diameter of Laval nozzle was 25.3 and 32.9 mm, respectively. There were three concentric rings supplying the shrouding CH_4 and O_2 , as presented in Fig. 1a. The inner hydraulic, intermediate and outer hydraulic diameter is 3.2, 4.5 and 4.8 mm, respectively. For combustion experiment and numerical simulation, the same parameters of supersonic jet nozzle were applied in this paper.

The point tube with water-cooling was adopted to measure the static pressure and total pressure of coherent jet in the combustion experiment, as shown in the Fig. 1b. The Mach number of coherent jet should be calculated by following formation (Anderson 2013):

$$Ma^2 = \frac{2}{\gamma - 1} \left[\left(\frac{p_0}{p} \right)^{(\gamma-1)/\gamma} - 1 \right] \quad (1)$$



where P_0 is the pressure at the A location, P is the pressure at the B location and γ is the ratio of heat capacity.

Numerical model

Governing equations

The numerical simulations were conducted by integrating the Navier–Stokes equations with Reynolds averaging method, and following equations was used to simulated combustion process (Malalasekera and Versteeg 2007).

Mass conservation equation

$$\frac{\partial \rho}{\partial \tau} + \frac{\partial \rho U_i}{\partial X_i} = 0 \tag{2}$$

Momentum conservation equation

$$\frac{\partial}{\partial \tau}(\rho \vec{v}) + \nabla \cdot (\rho \vec{v} \vec{v}) = -\nabla P + \nabla \left[\mu_e (\nabla \vec{v} + \nabla \vec{v}^T) \right] + \rho \vec{g} + f_\alpha \tag{3}$$

Energy equation

$$\frac{\partial}{\partial \tau}(\rho E) + \nabla \cdot (\vec{v} (\rho E + \rho)) = \nabla (k_{eff} \nabla T - \sum_j h_i \vec{J}_j + (\vec{\tau}_{eff} \vec{v})) + S_h \tag{4}$$

In (2–4), \vec{v} was the velocity vector; k_{eff} was the effective conductivity and \vec{J}_j was the diffusion flux of species j ; S_h included the heat of chemical reaction, and any other volumetric heat sources.

The discrete ordinate (DO) radiation model (Christo and Dally 2005; Chui and Raith 1993) with weighted sum of gray gas model (WSGGM) was used to model the radiation of the combustion process. In the present study, the modified $k-\epsilon$ model with the standard wall function was implemented for modeling the turbulent flows. The standard $k-\epsilon$ model (Launder and Spalding 1972) was a semi-empirical model based on model transport equations for the turbulence kinetic energy (k) and its dissipation rate (ϵ).

Turbulence kinetic energy equation for EDC model (k equation):

$$\frac{\partial(\rho k)}{\partial \tau} + \frac{\partial}{\partial x_i}(\rho k u_i) = \frac{\partial}{\partial x_j} \left[\left(\mu + \frac{\mu_t}{\sigma_k} \right) \frac{\partial k}{\partial x_j} \right] + G_k + G_b - \rho \varepsilon - Y_M + S_k \quad (5)$$

Turbulence dissipation equation for EDC model (ε equation):

$$\frac{\partial(\rho \varepsilon)}{\partial \tau} + \frac{\partial}{\partial x_i}(\rho \varepsilon u_i) = \frac{\partial}{\partial x_j} \left[\left(\mu + \frac{\mu_t}{\sigma_\varepsilon} \right) \frac{\partial \varepsilon}{\partial x_j} \right] + C_{1\varepsilon} \frac{\varepsilon}{\kappa} (G_k + C_{3\varepsilon} G_b) - C_{2\varepsilon} \rho \frac{\varepsilon^2}{\kappa} + S_\varepsilon \quad (6)$$

In the equations, the G_k and G_b presented was the generation of turbulence kinetic energy due to the mean velocity gradient and the buoyancy, respectively. Y_M represented the contribution of the fluctuating dilatation in compressible turbulence to the overall dissipation rate. S_k and S_ε were user-defined source terms. In combustion process, the temperature gradient was an important factor to obtain the flow field (Mardani et al. 2010; Magnussen and Hjertager 1977). Hence, the turbulent viscosity (μ_t) is addressed as the following equation (Alam et al. 2010b):

$$\mu_t = \frac{0.09 \rho \varepsilon^2}{\left[1 + \frac{1.2 T^{0.6}}{1+f(M_\tau)} \right] k} \quad (7)$$

The EDC model with overall and detailed chemical kinetic mechanisms (GRI-Mech 3.0) were presently used for the modeling of reactions (Frassoldati et al. 2010; Mardani and Tabejama 2010; Galletti et al. 2009). The GRI-Mech 3.0 was the full chemical mechanism, which consisted of 53 species and 325 reversible reactions. In EDC model, the species conservation equation for chemical species taken the following general form:

$$\frac{\partial}{\partial t}(\rho Y_i) + \nabla \cdot (\rho \vec{v} Y_i) = -\nabla \cdot \vec{J}_i + R_i \quad (8)$$

where Y_i was the local mass fraction of each species (i), \vec{J}_i was the diffusion flux of species, and R_i was the net rate of production of species by chemical reaction. The length fraction of the fine scales (ξ) and the residence chemical time scale (τ) of fluid in the fine structures was expressed by:

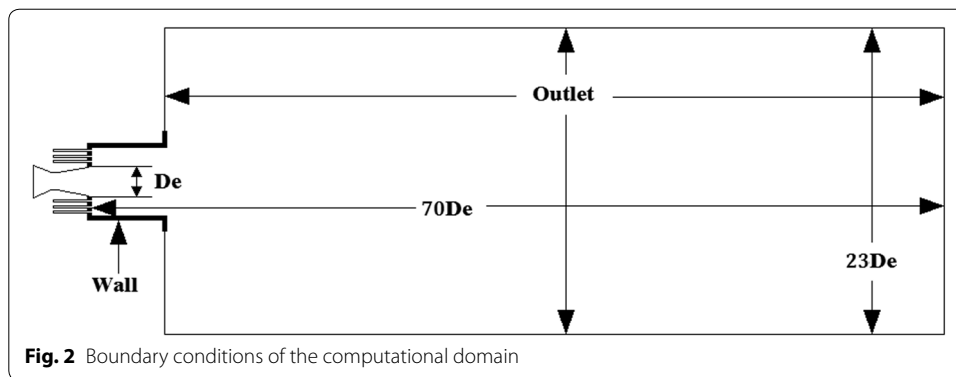
$$\xi = C_\xi \left(\frac{\nu \varepsilon}{k^2} \right)^2, \quad \tau = C_\tau \left(\frac{\nu}{\varepsilon} \right)^{\frac{1}{2}} \quad (9)$$

where C_ξ and C_τ were time scale constants equal to 2.138 and 0.408, respectively.

Simulation details

Because the vortices in flow field could not be correctly reflected by the symmetry of the system, a 3D geometrical model of computation was constructed in Fig. 2. The computational domain started at the entrance of the coherent jet nozzle, and extended $70D_e$ downstream in the axial direction and $23D_e$ in the radial direction.

The main oxygen used pressure inlet condition, surrounding CH_4 and O_2 adopted mass-flow inlet condition for and outlet position of combustion zone adopted pressure-outlet. The Table 1 listed the detail values of boundary conditions.

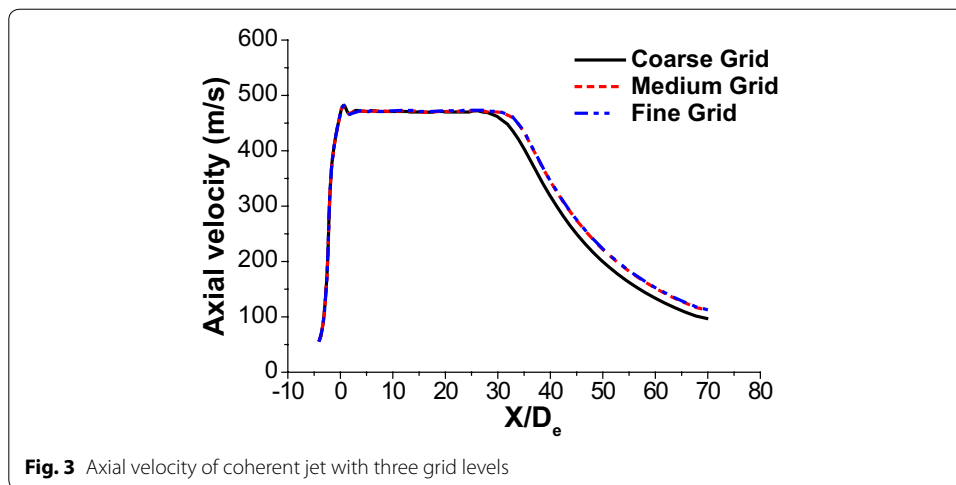
**Table 1** Specifications of boundary conditions

Name of boundary	Type of boundary conditions	Values
Main oxygen inlet	Stagnation pressure (Pa)	768,545
	Mach number	2.0
	Mass fractions (wt%)	O ₂ = 100
	Oxygen temperature (K)	298
CH ₄ inlet	Mass flow rate (kg/s)	0.0456/0.0411/0.0365/0.0319/0.0274/0.0228/0
	Mass fractions (wt%)	CH ₄ = 100
	CH ₄ temperature (K)	298
Surrounding oxygen inlet	Mass flow rate (kg/s)	0.0913/0.0821/0.0730/0.0639/0.0548/0.0456/0
	Mass fractions (wt%)	O ₂ = 100
	Surrounding oxygen temperature (K)	298
	Outlet	Static pressure (Pa)
	Mass fractions (wt%)	O ₂ = 21, N ₂ = 79
Ambient temperature	Initial temperature (K)	298/1700

In present study, The SIMPLE algorithm method (Chui and Raith 1993; Launder and Spalding 1972) was utilized to solve pressure velocity coupling. The second-order upwind scheme was employed for discretizing the equations in order to improve the accuracy of the simulations. Solution convergence was determined by two criteria. The first one was to ensure that the numerical residuals were $<10^{-6}$ for the energy and $<10^{-5}$ for all the other variables. The second criterion was to ensure that the variations between consecutive iterations of temperature and velocity were kept within 10 K and 2 m/s, respectively, at the downstream outlet of the computational domain.

There were large quantities of previous studies (Mardani et al. 2010; Magnussen and Hjertager 1977; Frassoldati et al. 2010; Mardani and Tabejama 2010; Galletti et al. 2009) had used GRI-Mech 3.0 to model the combustion process. Thus, the EDC model with the detailed chemical kinetic mechanism of GRI-Mech 3.0 was used for the modeling of reactions. To reduce the calculation time, the in situ adaptive tabulation (ISAT) model of Pope (1997) was adopted.

Figure 3 shows the axial velocity profiles for the simulation results in three grid levels, coarse grid (275,000 cells), medium grid (417,000 cells), and fine grid (564,000 cells), respectively. It seems that the simulation results agree well with the each other, for the



medium and the fine grid, and the variation of axial velocity is within 0.7 pct. This result suggests that the solution is not sensitive to the grid. However, the variation calculated with the coarse and medium grid level was about 2.4 pct. Therefore, the mesh of 417,000 cells was used in combustion simulation in consideration of reducing the computational time.

Results and discussions

Axial velocity distribution research

Figure 4 shows the axial velocity distribution of the supersonic oxygen jet with different shrouding CH₄ flow rates and ambient temperatures. Moreover, the simulation results would be addressed as line or line segment, and the experiment data would be presented as quadrangle, roundness or triangle in Figs. 4 and 6. The mass flow rates of CH₄ in different unit are shown in the Table 2. In this paper, the percent of main oxygen design flow rate will be used to replace CH₄ mass flow rate, hereafter. For instance, 10 %-CH₄ represents the shrouding CH₄ mass flow rate being 0.0456 kg/s.

Based on average values of different kinds of working conditions in Fig. 4, within 5 nozzle exit diameters from the nozzle exit, all the velocities of main oxygen jets are

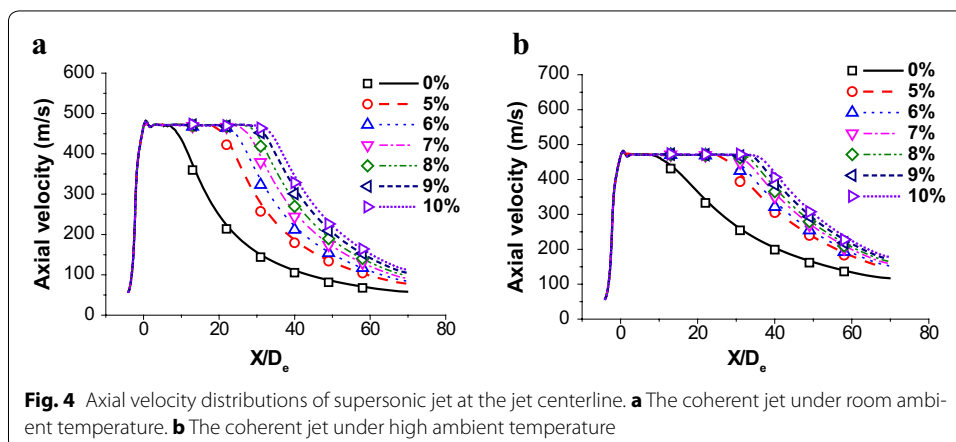


Table 2 The mass flow rates of CH₄ in different unit

Mass flow rate of CH ₄ (kg/s)	Pneumatic volume flow rate of CH ₄ (Nm ³ /h)	Pneumatic volume flow rate of O ₂ (Nm ³ /h)	The percent of main oxygen flow rate (%)
0.0456	230	2300	10
0.0411	207		9
0.0365	184		8
0.0319	161		7
0.0274	138		6
0.0228	115		5
0	0		0

471 ± 10 m/s at the jet centerline. As the results show, more shrouding CH₄ rate does not increase the maximum velocity of main oxygen jet, and it just makes the maximum velocity remain a longer distance at the jet centerline. Therefore, CH₄ flow rate could only protect the stirring ability of main oxygen jet, but not improve the stirring ability at a certain extent.

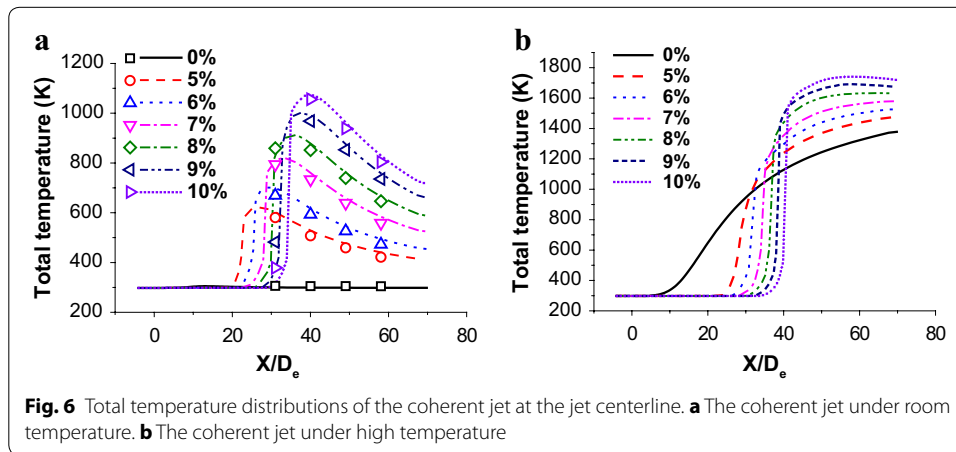
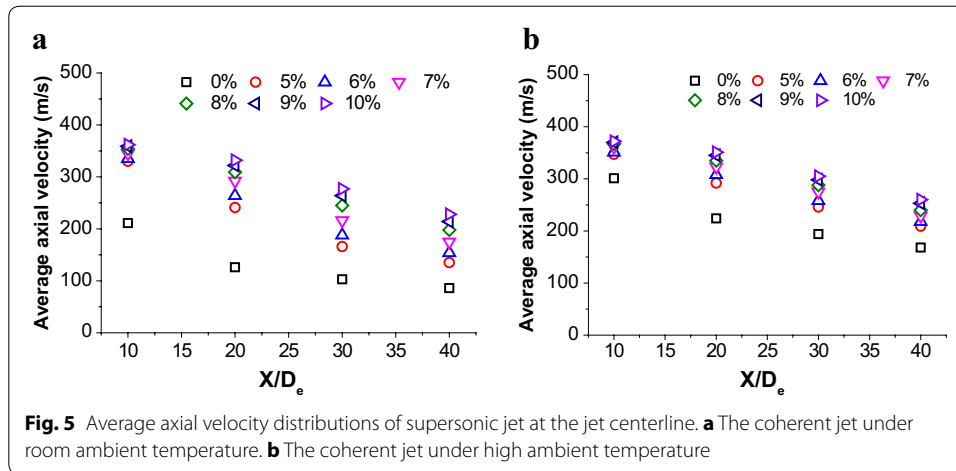
1 nozzle exit diameter will be addressed as 1D_e hereafter. The potential core of coherent jet is supposed to be end, when the axial velocity of supersonic oxygen jet continues to decline more than 5D_e at the jet centerline, and the first point of this kind of down-trend is addressed as the end point of potential core. The length between the end point of potential core and the tip of Laval nozzle is defined as the potential core length of coherent jet.

As shown in Fig. 4a, the potential core length with 10 %-CH₄, 9 %-CH₄, 8 %-CH₄, 7 %-CH₄, 6 %-CH₄ and 5 %-CH₄ is 29D_e, 27D_e, 25D_e, 23D_e, 21D_e and 18D_e, respectively, at room ambient temperature. Under high ambient temperature, the potential core length with 10 %-CH₄, 9 %-CH₄, 8 %-CH₄, 7 %-CH₄, 6 %-CH₄ and 5 %-CH₄ is 34D_e, 32D_e, 30D_e, 28D_e, 26D_e and 24D_e, respectively, as depicted in Fig. 4b. As previous studies (Mathur 1999b), the higher ambient temperature can prolong potential core of coherent jet, and the more CH₄-rate also could increase the potential core referring the results.

When the ambient temperature being 1700 K, the length of potential core with 10 %-CH₄, 9 %-CH₄, 8 %-CH₄, 7 %-CH₄, 6 %-CH₄ and 5 %-CH₄ is 1.17, 1.18, 1.20, 1.22, 1.24 and 1.33 times larger than that of the coherent jet at ambient temperature being 298 K, respectively. Therefore, higher ambient temperature would prolong the potential core length, but the increasing trend decelerates with shrouding rate increasing.

Based on the results, the average potential core length of the coherent jet is 3.5 times larger than that of the conventional jet. That means the CH₄ flow rate could effectively improve the potential core length of the main oxygen jet, which protect the stirring ability of the supersonic oxygen jet.

There are four cross sections have been selected at the jet centerline which radius is 100 mm, and the location is 10D_e, 20D_e, 30D_e and 40D_e, respectively. The Fig. 5 shows the average axial velocities of different sections, which are calculated by numerical simulation. Although the coherent jet have a potential core at centerline, the average velocity shows a reduce trend, because periphery of main oxygen would mix with the combustion flame, which suppresses the average velocity of oxygen jet.



The average velocity of the coherent jet at high temperature condition with 10 %-CH₄, 9 %-CH₄, 8 %-CH₄, 7 %-CH₄, 6 %-CH₄ and 5 %-CH₄ is 1.07, 1.09, 1.11, 1.16, 1.21 and 1.25 times higher than that of the coherent jet at room ambient temperature, respectively. Therefore, if the ambient temperature rises, the increasing trend of average velocity is suppressed with CH₄ flow rate improving, which is the same as increasing trend of potential core.

In this research, the distance between the coherent jet tip and molten bath surface is approximately 20D_e. When the location of cross section is 20D_e at centerline, the average velocity of the coherent jet at high temperature condition with 9 %-CH₄, 8 %-CH₄, 7 %-CH₄, 6 %-CH₄ and 5 %-CH₄ is 98.3, 95.4, 92.0, 87.7 and 83.2 % of the 10 %-CH₄ coherent jet, respectively. It seems that when the CH₄ flow rate is 9 %-CH₄ and 8 %-CH₄, the disparity rate of average velocity is all small than 5 %.

Based on the preview research (Anderson 2013; Alam et al. 2012; Hale 2013), the velocity contour of main oxygen jet would be compressed due to the obstruction of molten bath surface. Therefore, the difference of main oxygen velocity with 10 %-CH₄, 9 %-CH₄ and 8 %-CH₄ may be suppressed at 20D_e, which makes the coherent jet may

achieve the same stirring effect with three kinds of CH₄ flow rate in EAF steelmaking process.

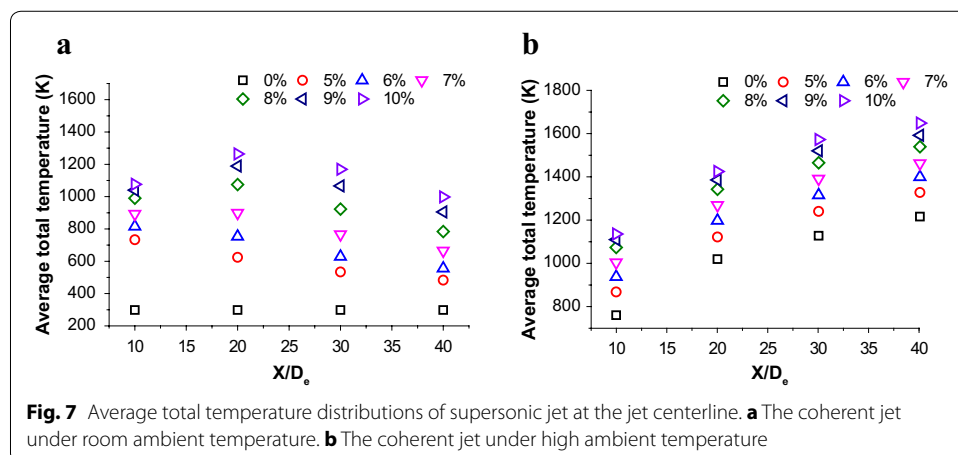
The Fig. 4 presents numerical simulation results are agreed well with the combustion experiment in this research, and the average error of axial velocity is about 1.7 %. The oxygen is addressed as an ideal-gas in simulation process. However, the oxygen has a deviation from the ideal gas, when the ambient temperature is improving, referring to the previous study (Irvin and Richard 2008). Therefore, the error would be formed under the influence of shrouding flame temperature.

Total temperature distribution research

As depicted in Fig. 6, the simulations using the EDC model with full detailed chemistry are performed to predict the flame characteristics with various CH₄ flow rates. In this study, the total temperature distributions at the jet centerline are measured only at cold ambient temperature, and the simulation results agree well with the combustion experiment data.

When oxygen jet passes though the Laval nozzle, the pressure potential energy is transformed into the kinetic energy. During this process, although the energy form of oxygen jet is changing, the total energy of oxygen remains unchanged, and then the central jet mixes with combustion flame at the end of the potential core. With a great temperature gradient, the flame transmits thermal energy into oxygen jet, which makes centre of oxygen jet rise rapidly. At last, both kinetic energy and thermal energy of oxygen jet keeps reducing due to the energy gradient between the jet and the ambient flow. As a result, total temperature of the jet gradually approaches the ambient temperature. As described above, the total temperature at jet centerline just one of flow field characteristics, which could not represent the total stirring ability of coherent jet. The average total temperatures of different sections are studied, as shown in Fig. 7.

Under room ambient temperature, when the CH₄ flow rate is small than 6 %-CH₄, the average total temperature show a reduce trend with X/D_e increasing. However, when the CH₄ flow rate is bigger than 7 %-CH₄, the average total temperature increases first, and then approaches to the ambient temperature at centerline direction, as presented in Fig. 7a.



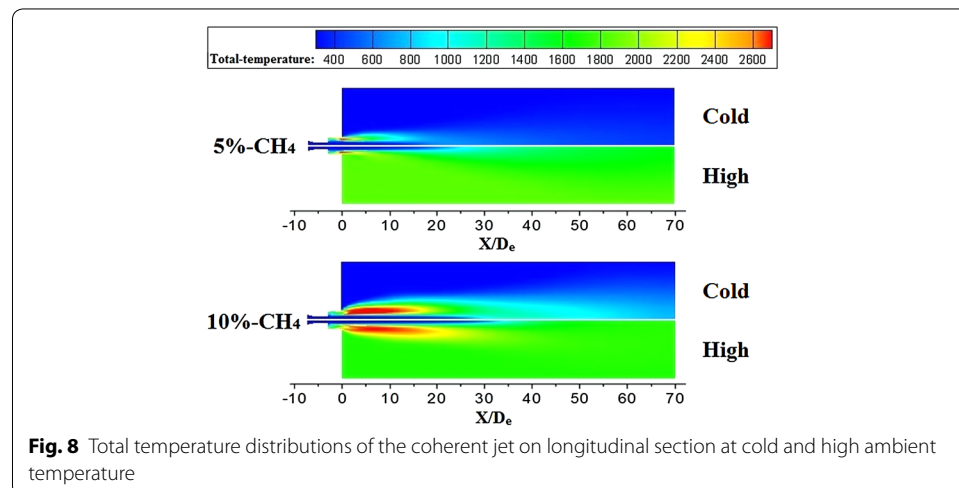
It seems that the high-temperature zone (the total temperature >2600 K) is formed near the exit of coherent, when CH_4 flow rate is small, as depicted in Fig. 8. With the combustion flame flowing to the downstream, the thermal energy of coherent jet transmits into the ambient gas due to the temperature gradient, which makes its average total temperature keep reducing. When CH_4 flow rate is bigger, the high-temperature zone is enlarged making the average total temperature of combustion flame higher at $20X/D_e$. At the same time, the average total temperature of main oxygen increase, because the peripheral of oxygen jet absorbs heat from flame. As a result, the average total temperature at $20X/D_e$ is higher than that at $10X/D_e$. And then, the total temperature reduce rate of flame is bigger than total temperature rise rate of oxygen jet, which makes the average total temperature of coherent jet decreases with X/D_e increasing.

When the ambient temperature is high, the average total temperature shows an increasing trend with different CH_4 flow rates, as presented in Fig. 7b. Although the CH_4 has been exhausted, the combustion flame could keep its total temperature being 1700 K, because of the high ambient temperature. In the meantime, the main oxygen jet could keep absorbing heat until its total temperature being 1700 K. Therefore, on the cross sections, the average total temperature of coherent jet increases, because of the limited reduction rate of flame total temperature and increasing trend of oxygen jet total temperature.

The average total temperature of the coherent jet at high temperature condition with 10 %- CH_4 , 9 %- CH_4 , 8 %- CH_4 , 7 %- CH_4 , 6 %- CH_4 and 5 %- CH_4 is 1.28, 1.33, 1.43, 1.58, 1.76 and 1.92 times higher than that of the coherent jet at room ambient temperature, respectively. Therefore, if the ambient temperature rises, the increasing trend of average total temperature is suppressed with CH_4 flow rate improving, which is the same as increasing trend of average velocity.

When the location of cross section is $20D_e$ at centerline, the average total temperature of the coherent jet at high temperature condition with 9 %- CH_4 , 8 %- CH_4 , 7 %- CH_4 , 6 %- CH_4 and 5 %- CH_4 is 97.3, 94.2, 89.1, 84.1 and 78.8 % of the 10 %- CH_4 coherent jet, respectively.

It seems that when the CH_4 flow rate is 9 %- CH_4 and 8 %- CH_4 , the disparity rate of average velocity is all small than 6 %. Therefore, the 9 %- CH_4 and 8 %- CH_4 may achieve the same stirring effect in EAF steelmaking process.



Industrial application research

Based on the combustion experiment and numerical simulation study, the 10 %-CH₄, 9 %-CH₄ and 8 %-CH₄ may achieve the same stirring effect in a 75t electrical arc furnace. In order to study metallurgical effects and technical indicators of the various CH₄ flow rate in steelmaking process, 10 %-CH₄, 8 %-CH₄ and 6 %-CH₄ are adopted in a 75 t electrical arc furnace. There are 180 heats collected in the industrial smelting process, and each CH₄ flow rate has same heats. Molten steel component, steelmaking time and dephosphorization rate are analyzed in this research.

The conditions of liquid iron (prior to steelmaking process), and molten steel (after steelmaking process) are shown in Table 3 along with the average components, temperature, and smelting time. When smelting with different CH₄ flow rates, the conditions of liquid iron are fundamentally the same, which mean the initial conditions have no influence to the industrial application research.

Based on the carbon content and temperature of molten steel are same, both 10 %-CH₄ and 8 %-CH₄ make no difference on [C]·[O] and steelmaking time. When the CH₄ flow rate is 6 %-CH₄, the average [C]·[O] is improved by 16.2 % and the steelmaking time is increased by 7.4 %. This fact can prove that the stirring ability of coherent jet with 10 %-CH₄ and 8 %-CH₄ is fundamentally same. Therefore, the dynamic condition of molten bath remains unchanged, when the CH₄ flow rate is reduced at appropriate situation as observed by the combustion experiment and numerical simulation.

Figure 9 shows the distribution of phosphorus in molten steel with different CH₄ flow rates. As shown in Fig. 9, the content of phosphorus in molten steel with 10 %-CH₄, 8 %-CH₄ and 6 %-CH₄ distributes from 0.006 to 0.008, 0.006 to 0.008 and 0.009 to 0.011 mass%, respectively. Based on the results, the average content of phosphorus in molten steel is basically same when the CH₄ flow rate is 10 %-CH₄ and 8 %-CH₄. However, the average content of phosphorus is increased by 0.003 mass % with 6 %-CH₄, and the dephosphorization rate is reduced by 2.3 %.

CaO, SiO₂, FeO and P contents in end-point slag are shown in Fig. 10. While the basicities of slag are same being 2.1, compared with 6 %-CH₄, average P is increased by 0.07 mass% and FeO is dropped by 2.2 mass% in slag with 10 %-CH₄ and 8 %-CH₄. The reduction of FeO loss will be beneficial to improve the metal yield rate.

Based on steelmaking temperature, components of slag layer and molten steel, the equilibrium phosphorus content with various CH₄ flow rate has been obtained by the Eq. 10 reported by Healy (1970).

$$\log \frac{(\text{mass\% P})}{[\text{mass\% P}]} = \frac{22,350}{T} + 0.08(\text{mass\% CaO}) + 2.5 \log(\text{mass\% FeO}) - 16 \quad (10)$$

Table 3 Average values analysis of liquid iron and molten steel

Label	Liquid iron			Molten steel				Steelmaking time (min)
	C (%)	P (%)	Temperature (°C)	C (%)	P (%)	[C]·[O] (10 ⁻⁴)	Temperature (°C)	
10 %-CH ₄	3.6	0.129	1280	0.078	0.007	0.0053	1605	50.3
8 %-CH ₄	3.6	0.130	1279	0.079	0.007	0.0052	1606	50.4
6 %-CH ₄	3.6	0.129	1280	0.081	0.010	0.0061	1605	54.1

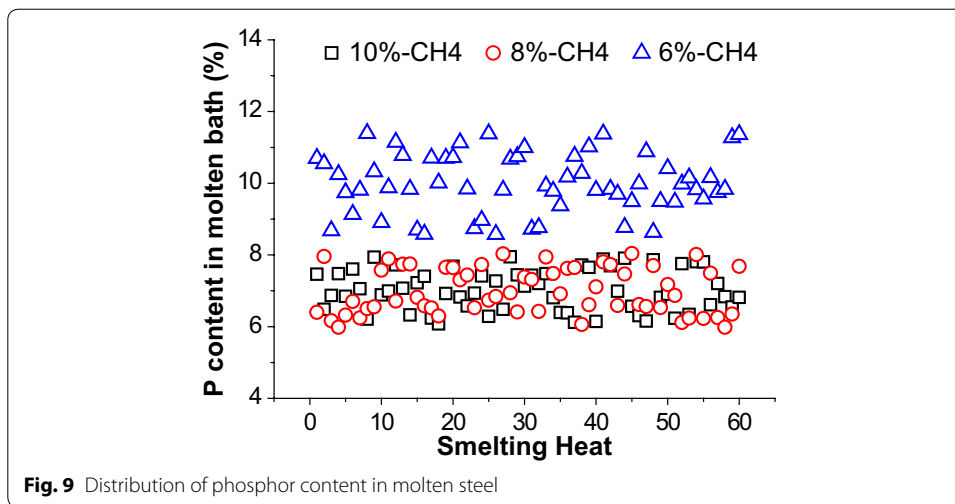


Fig. 9 Distribution of phosphor content in molten steel

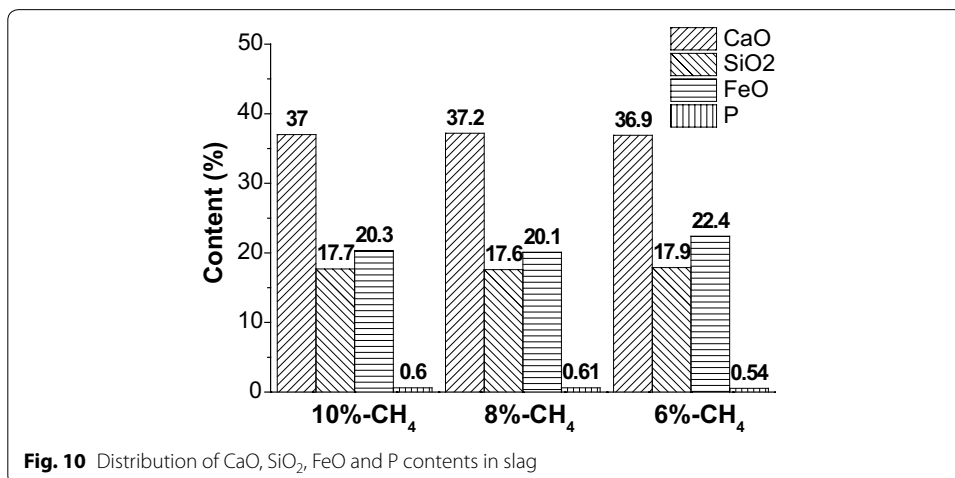


Fig. 10 Distribution of CaO, SiO₂, FeO and P contents in slag

The content of equilibrium phosphorus in molten steel is 4.2×10^{-5} , 4.3×10^{-5} and 3.0×10^{-5} mass% with 10 %-CH₄, 8 %-CH₄ and 6 %-CH₄, respectively. It is obvious that all equilibrium phosphorus contents are most importantly negligibly small. According to the results of combustion experiment, numerical simulation and industrial application research, it can prove that the reaction rate and stirring effect with 10 %-CH₄ and 8 %-CH₄ is same, and both are better than that with 6 %-CH₄.

Above all, it shows that the stirring ability of coherent jet with 10 %-CH₄ and 8 %-CH₄ is fundamental same, and both metallurgical effect is better than with 6 %-CH₄ in steel-making process for the 75 t EAF, which agrees well with the results of the combustion experiment and the numerical simulation.

Conclusions

1. The bigger CH₄-rate and higher ambient temperature can prolong the potential core of coherent jet, and the numerical simulation results show a good agreement with the combustion experimental data. Moreover, when the ambient temperature rises,

the increasing trend of potential core is suppressed with CH_4 flow rate improving, which is the same as increasing trend of average velocity and total temperature.

- Based on the results, at high temperature condition, the disparity rate of average velocity and total temperature of the coherent jet, which is between 10 %- CH_4 , 9 %- CH_4 and 8 %- CH_4 , is all small than 5 and 6 %, respectively. Therefore, the coherent jet may achieve the same stirring effect with three kinds of CH_4 flow rates in EAF steelmaking process.
- Compared with 10 %- CH_4 and 6 %- CH_4 flow rates at industrial application research, the best CH_4 flow rate is 8 %- CH_4 , which could stir molten bath well and reduce energy consumption in steelmaking process, with the condition of liquid iron being fundamentally the same.

Authors' contributions

All authors participated in the preparation of the manuscript. All authors read and approved the final manuscript.

Author details

¹ National Center for Materials Service Safety, University of Science and Technology Beijing, Beijing 100083, China. ² Beijing Key Laboratory of Research Center of Special Melting and Preparation of High-end Metal Materials, University of Science and Technology Beijing, Beijing 100083, China. ³ Rongcheng Jingye Technology Limited Company Beijing, Beijing 100083, China. ⁴ Tianjin Pipe Co. Ltd., Tianjin 300301, China.

Competing interests

The authors declare that they have no competing interests.

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References

- Alam M, Naser J, Brooks G, Fontana A (2010a) Computational fluid dynamics modeling of supersonic coherent jets for electric arc furnace steelmaking process. *Metall Trans B* 41:1354
- Alam M, Naser J, Brooks G, Fontana A (2010b) Computational fluid dynamics simulation of supersonic oxygen jet behavior at steelmaking temperature. *Metall Trans B* 41B:636
- Alam M, Naser J, Brooks G, Fontana A (2012) A computational fluid dynamics model of shrouded supersonic jet impingement on a water surface. *ISIJ Int* 52:1026
- Anderson JD (2013) Introduction to flight. McGraw-Hill Education (Asia), Singapore, p 125
- Anderson JE, Mathur PC, Selines RJ (1998) Method for introducing gas into a liquid. US Patent 5814125
- Christo FC, Dally BB (2005) Modeling turbulent reacting jets issuing into a hot and diluted co-flow. *Combust Flame* 142:117
- Chui EH, Raith GD (1993) Computation of radiant heat transfer in a nonorthogonal mesh using the finite-volume method. *Numer Heat Transf B Fund* 23:269
- Deo B, Boom R (1993) Fundamentals of steelmaking metallurgy. Prentice Hall, Upper Saddle River
- Frassoldati A, Sharma P, Cuoci A, Faravelli T, Ranzi E (2010) Kinetic and fluid dynamics modeling of methane/hydrogen jet flames in diluted coflow. *Appl Therm Eng* 30:376
- Galletti C, Parente A, Derudi M, Rota R, Tognotti L (2009) Numerical and experimental analysis of NO emissions from a lab-scale burner fed with hydrogen-enriched fuels and operating in MILD combustion. *Int J Hydrogen Energy* 34:8339
- Hale FJ (2013) Introduction to aircraft performance, selection, and design. Wiley, New York, p 1
- Healy GW (1970) A new look at phosphorous distribution. *J Iron Steel Inst* 208:664
- Irvin G, Richard AY (2008) Combustion. Academic Press, Burlington
- Jeong MS, Kumar VRS, Kim H-D, Setoguchi T, Matsuo S (2004) A computational characterization of the supersonic coherent jet. In: 40th AIAA/ASME/SAE/ASEE joint propulsion conference, Fort Lauderdale, FL
- Launder BE, Spalding DB (1972) Lectures in mathematical model of turbulence. Academic Press, London, p 124
- Liu J, Warner AEM, McCann D, Hall DE, Mallette D, Bradley JA, Mackenzie E, Mahoney WJ, Deneys A (2005) Converter and fire refining practices. TMS, San Francisco, p 61
- Liu F, Zhu R, Dong K, Hu S (2016) Effect of ambient and oxygen temperature on flow field characteristics of coherent jet. *Metall Trans B* 47B:228
- Magnussen BF, Hjertager BH (1977) On mathematical models of turbulent combustion with special emphasis on soot formation and combustion. *Proc Combust Inst* 16:719

- Mahoney WJ (2008) High mach number, sub-atmospheric, coherent jets and their application in vacuum metallurgy. In: 3rd international conference on process development in iron and steelmaking, vol 1, p 367
- Mahoney WJ (2010) Experimental remarks on supersonic jet behavior in high temperature, reactive ambient in connection to steelmaking. In: AISTech-iron and steel technology conference proceedings, Pittsburgh, PA, p 1071
- Malalasekera W, Versteeg HK (2007) An introduction to computational fluid dynamics. Harlow, Pearson, p 1
- Mardani A, Tabejama S (2010) Effect of hydrogen on hydrogen–methane turbulent non-premixed flame under MILD condition. *Int J Hydrogen Energy* 35:11324
- Mardani A, Tabejamaat S, Ghamari M (2010) Numerical study of influence of molecular diffusion in the mild combustion regime. *Combust Theory Model* 14:747
- Mathur PC (1999a) Fundamentals and operating results of praxair's Cojet technology. *Iron Steelmaker* 26:59
- Mathur PC (1999b) Coherent jets in steelmaking: principles and learnings. *Iron Steelmak* 3:50
- Meidani ARN, Isac M, Richardson A, Cameron A, Guthrie RIL (2004) Modelling shrouded supersonic jets in metallurgical reactor vessels. *ISIJ Int* 44:1639
- Naito K, Ogawa Y, Inomoto T, Kitamura S, Yano M (2000) Characteristics of jets from top-blown lance in converter. *ISIJ Int* 40:23
- Nordquist A, Kumbhat N, Jonsson L, Jönsson P (2006) The effect of nozzle diameter, lance height and flow rate on penetration depth in a top-blown water model. *Steel Res Int.* 77:82
- Pope SB (1997) Computationally efficient implementation of combustion chemistry using in situ adaptive tabulation. *Combust Theory Model* 1:41
- Sarma B, Mathur PC, Selines RJ, Anderson JE (1998) Fundamental aspects of coherent gas jets. *Electric Furnace Conf Proc* 56:657
- Tago Y, Higuchi Y (2003) Fluid flow analysis of jets from nozzles in top blown process. *ISIJ Int* 43:209
- Wang W, Yuan Z, Matsuura H, Zhao H, Dai C, Tsukihashi F (2010) Three-dimensional compressible flow simulation of top-blown multiple jets in converter. *ISIJ Int* 50:491

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